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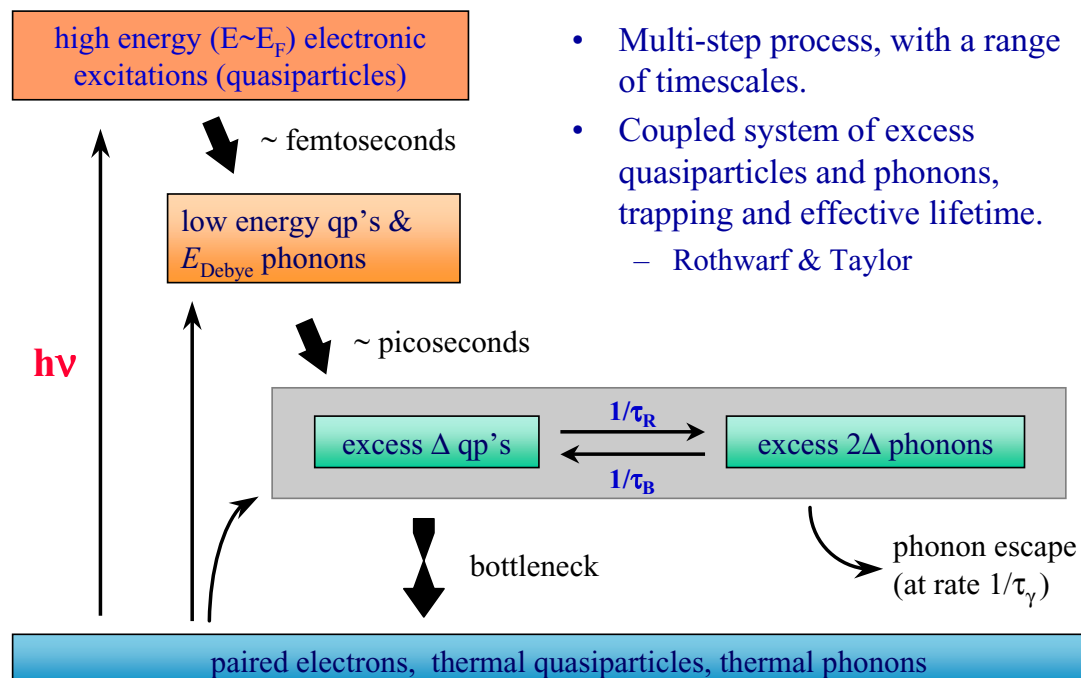
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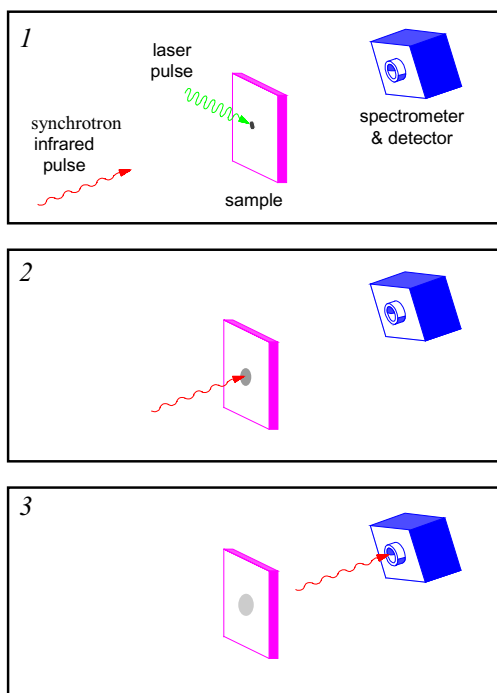
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OUTLINE

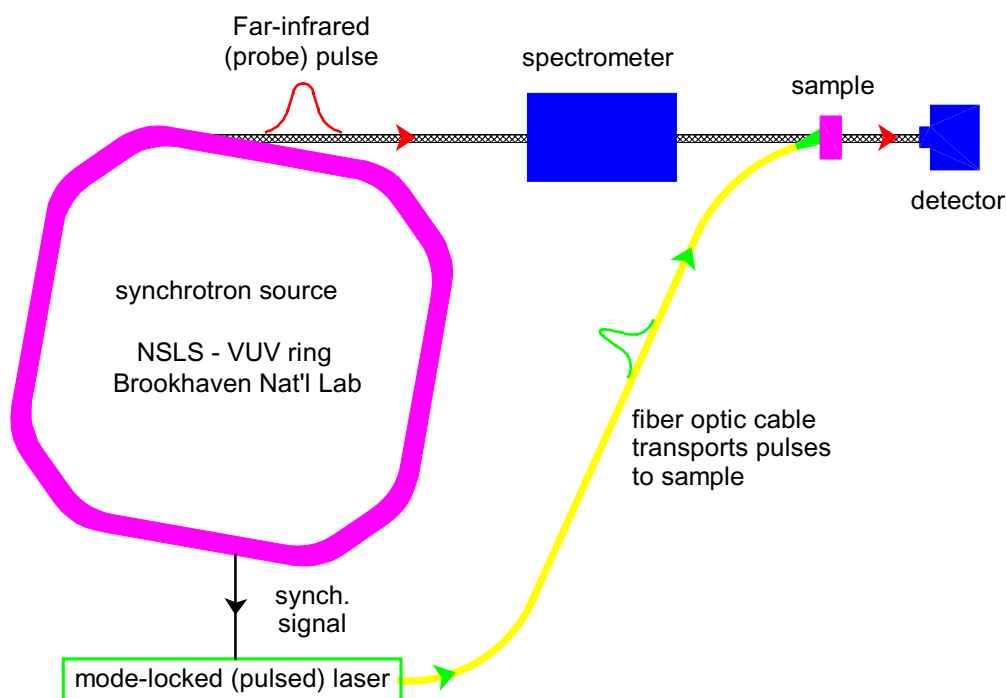
- Non-equilibrium state and relaxation process in superconductors
- Pump-probe technique
- MoGe
- Time-dependent relaxation of quasiparticles
- Conclusions



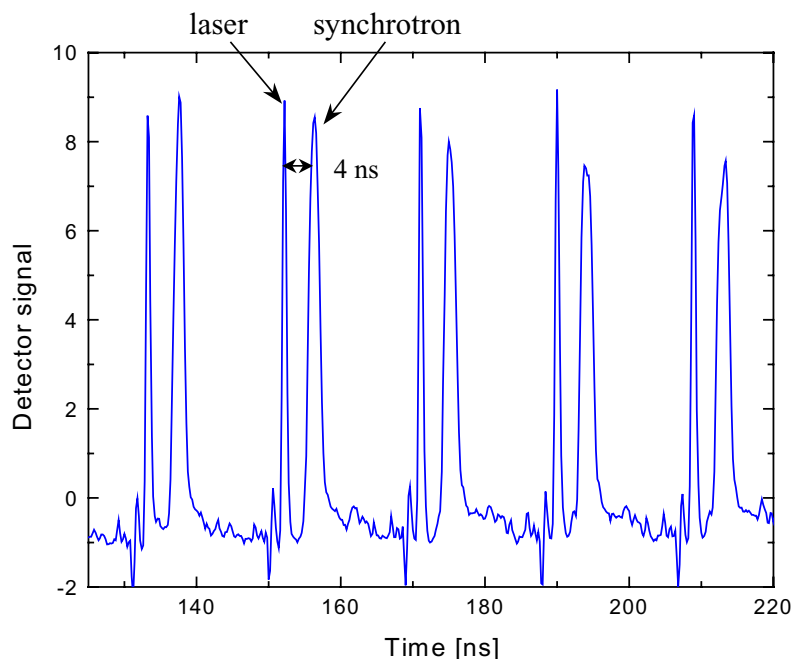


1. Laser pulse creates photoexcitations in sample, which subsequently evolve with time.
2. After time Δt , broadband (continuum) IR pulse arrives and is partially absorbed (or reflected) by excitations.
3. IR pulse analyzed with a spectrometer, extracting details of excitations at a time Δt after their creation.

- Cycle repeats at high (MHz) repetition rate.
- Photoexcitation evolution determined by measuring at a variety of Δt 's.
- Employs “conventional” spectroscopy using high-sensitivity (slow-response) detectors.



- Synchronized laser & storage ring pulses.
- Measured in sample chamber of U10A spectrometer.
- with Ge APD (near IR detector, ~1 ns response). Useful for locating “zero” delay point.
- Here, the delay is 4 ns.



FILMS:

Grown by RF sputtering
Ge buffer layer
Thickness measured with quartz thickness monitor.

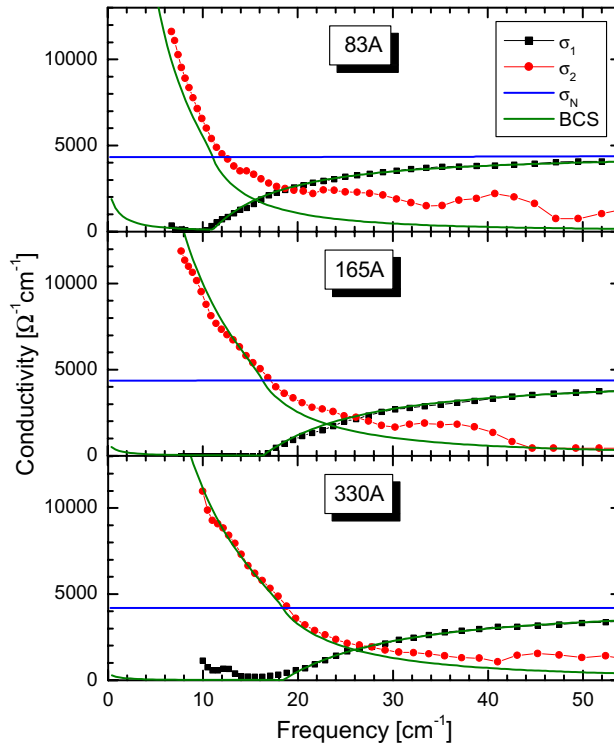
SUBSTRATE:

sapphire (r-cut)
1 mm thick
 $n_{\text{subs}} \sim 3.05$

KNOWN:

T_c bulk is about 7.2 K.
 T_c varies linearly with thickness and $1/R_{\square}$.
mean free path is of order 1-3 Angstroms

- A strong reduction of T_c with increasing R_{\square} is attributed to localization and related changes in the Coulomb interaction.
- α -MoGe serves as a model system for studying the interplay between superconductivity and disorder.
→ R_{\square} is the relevant measure of disorder in 2D.



Conductivity of a thin film on a thick substrate from transmittance and reflectance:

- Algorithm based on the approaches of Palmer and Tinkham and also Glover and Tinkham

R.E. Glover, III and M. Tinkham, *Phys. Rev.* 108, 243(1957)

L.H. Palmer and M. Tinkham, *Phys. Rev.* 165, 588(1968)

- σ_1 fits well to BCS for all three MoGe films.
- σ_2 has correct lineshape, $\sim 1/\omega$, but is above BCS, especially at higher frequencies.

Now define:

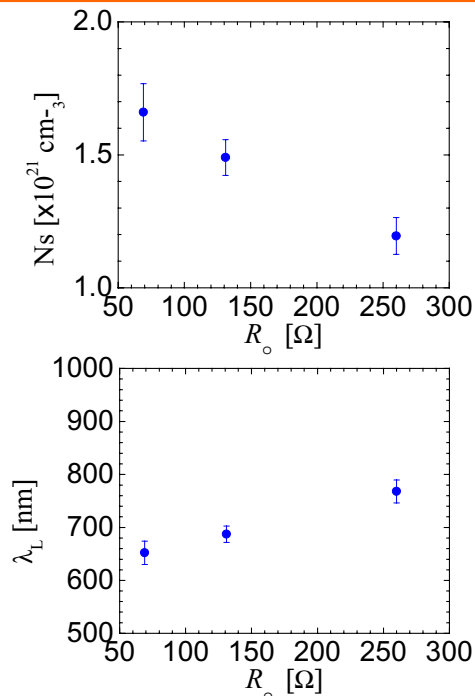
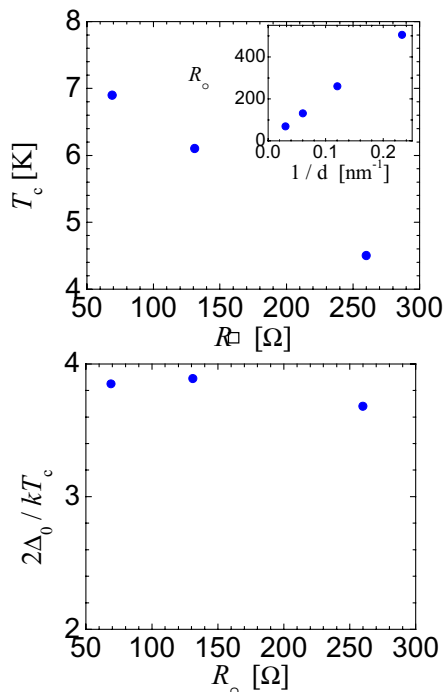
Superfluid density:

$$N_s(m/m^*) = mV_c\omega\sigma_2 / e^2$$

London penetration depth:

$$\lambda_L = c / [4\pi\omega\sigma_2]^{1/2}$$

How parameters depend on R_o



- T_c varies strongly with R_o and the superfluid density.
- $2\Delta_0 / kT_c$ does not depend strongly on R_o .

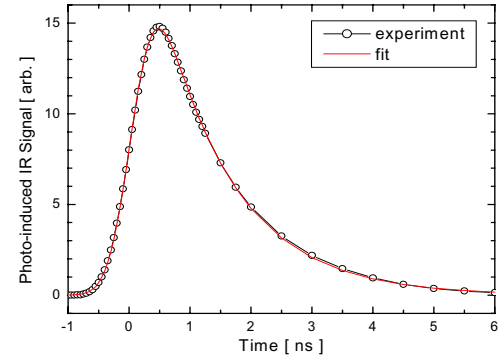
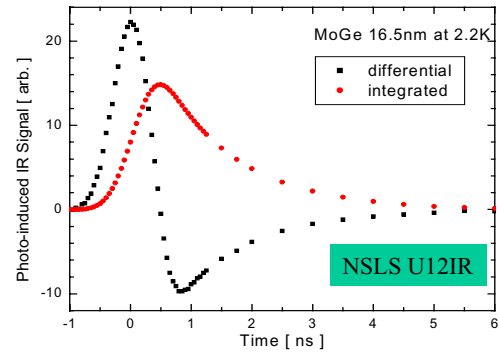
1 Differential Technique:

- pump-probe delay “dithered”.
- differential transmittance signal (spectral average) for a range of delay time.
- time-dependent relaxation of excess quasiparticles by integration.

1 Relaxation Behavior:

- convolution of simple exponential decay and Gaussian synchrotron pulse.
- decay time ~ 1 ns.
- time-resolution determined by synchrotron pulse width (> 300 ps).

$$\Delta T = \frac{1}{2} A \exp\left(\frac{w^2}{4\tau^2} - \frac{t-t_0}{\tau}\right) \left(1 + \operatorname{erf}\left(-\frac{w}{\tau} + \frac{t-t_0}{w}\right)\right)$$



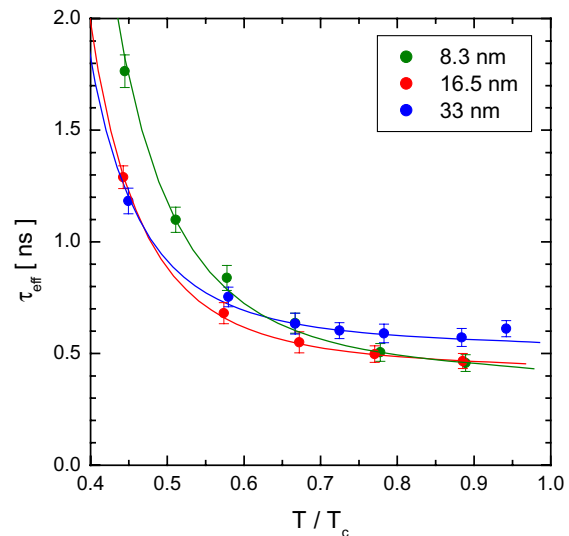
$$\tau_R(0) = (1 + \lambda) \hbar / 2\pi b (kT_c)^3$$

$$\tau_B(0) = \hbar N_\Omega / 4\pi^2 N(0) \langle \alpha^2 \rangle_{av} \Delta_0$$

$$\tau_{eff} \approx \tau_\gamma + (1/2) \tau_R (1 + \tau_\gamma / \tau_B)$$

- Relaxation times for MoGe films:

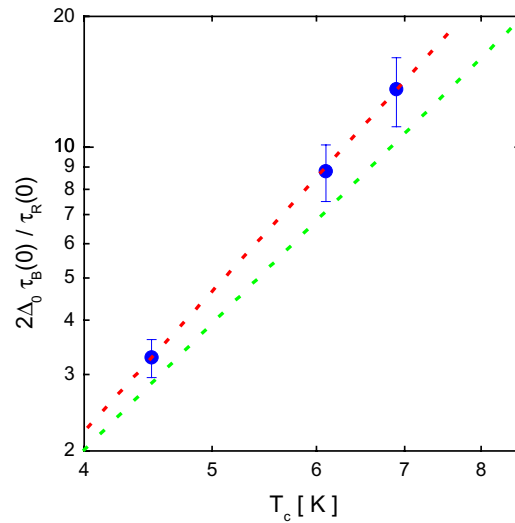
d [nm]	τ_γ [ps]	$\tau_R(0)$ [ps]	$\tau_B(0)$ [ps]	$\tau_R(0)/\tau_B(0)$
8.3	420	370	105	3.50
16.5	450	150	79	1.87
33	550	100	75	1.35



- Prediction by the theory (assuming that material parameters are the same for all three films):

$$\frac{2\Delta_0\tau_B(0)}{\tau_R(0)} \propto T_c^3$$

- Experiment:
 - The ratio $\sim T_c^3$.
 - Support the analysis of Kaplan et al.
 - Possible deviation due to invalid assumptions of the theory (e.g. weak-coupling, constant α^2 , simple Ω^2 dependence of phonon density of states, etc).



- Used pump-probe technique to follow relaxation of non-equilibrium superconducting a-MoGe films
- All three films fit dirty-limit BCS
- Our results are generally consistent with the theory by Kaplan et al.
- The effect of reduced thickness in these samples is to depress T_c and the superfluid density.
- At the same time, the normal-state conductivity and mean free path appear unchanged
- Indeed, it seems that, other than T_c , none of material parameters are changing with thickness.

- Faster FIR pulses – 1-10 ps
- Excellent FIR S/N $\sim 10^5/\text{rHr}$ (by power or stability or detector)
- Spectral coverage 1-100,000 cm^{-1}
- Modest resolution (0.5 cm^{-1} at 20 cm^{-1} is plenty)
- Low T capability: 1.7-300 K
- Range of pump wavelengths (FIR, MIR, NIR, VIS, UV)
- Range of pump fluences (up to $\sim 100 \text{ mW}$)
- Magnet?

(1) Ferromagnetic Semiconductors

- GaN:Mn, (ferromagnetic, Curie temperature $> 300\text{K}$)
- Pump-probe with a pump slightly larger than the GaN:Mn bandgap. (Need ~ 3 to $\sim 5 \text{ eV}$)

(2) Quantum confined Laser media

- III-V compounds.
- Carrier relaxation is critical to the lasing medium (determines population inversion)
- Pump at the bandgap, 800- 950 nm (Ti:sapphire) or 3-5 μm (OPA).
- Use THz as a probe of the carrier relaxation.

(3) Dual phase transitions: Multiferroics- Ferroelectric and Ferromagnetic

- Exhibit coexistence of two phase transitions: ferroelectricity + either antiferromagnetism or ferromagnetism.
- Example material: BiFeO_3 , Tb
- Pump carriers across the bandgap, using variable pump light polarization, and
- Study carrier thermalization and relaxation, in both the ordered and disordered states
- Need pump laser in the photon energy range of 3-5 eV.

(4) Cuprate superconductors

- Measure versus temperature and applied magnetic field,
- Measure transient change from photoexcitation.

(5) Ruthenates: Sr_2RuO_4 and homologous series

- Sr_2RuO_4 : T_c of 3 K (special case) or lower; 1.5 K more typical. Need a ^3He cooling system.
- Normal state differs qualitatively between Sr_2RuO_4 and SrRuO_3
- Sr_2RuO_4 has Fermi Liquid behavior over 1.5-30 K, crossover to non-Fermi liquid characteristics above $\sim 25\text{-}30\text{K}$.
- Measure THz absorption versus temperature

(6) Magnetic nanowire arrays

- Make magnetic nanowire arrays (Ni, Fe, Co, alloys) inside of alumina honeycomb arrays
- Wire diameter down to $\sim 5 \text{ nm}$ to $\sim 400 \text{ nm}$; 10-15% of the volume is wires
- Measure the THz absorption as a function of temperature and applied magnetic field
- Transient demagnetization studied by pump/probe (Ti:sapphire)

- Lighting consumes approximately 25% of all electrical power.
- Discharge sources typically 5 to 10x more efficient than incandescent sources
- High Intensity Discharge (HID) lamps, High Pressure Sodium (HPS) and Metal Halide (MH-HID), are in widespread use today.
- HID lamps operate in local thermodynamic equilibrium
- Total pressures: 1 to over 200 bar,
- Electron densities from 10^{15} - 10^{18} cm⁻³ (Plasma frequencies 0.3-10 THz)
- Diagnostic experiments on HID lamps are extremely challenging.
- Synchrotron radiation has the needed spectral radiance to "outshine" these very intense plasmas across broad spectral regions. (Optical and UV absorption at Stoughton / K-shell X-Ray at APS)
- Intense Terahertz radiation would be extremely valuable both for electron density measurements and for measuring the opacity of these lamps due to electron-atom inverse bremsstrahlung.
- Lamps are compact and well suited to experiments at remote sites.